

II, 1.6 Single-Particle Dynamics

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selected ('s') or tagged particle or
single ('s') particle, trajectory $\vec{R}(t)$

homogeneity in time: $\vec{R}(t-t_0)$ has properties
similar to $\vec{R}(t)$

homogeneity in space: $\vec{R}(t) + \vec{R}_0$ has similar
properties as $\vec{R}(t)$

stationary process $\Delta \vec{R}(t-t_0) = \vec{R}(t) - \vec{R}(t_0)$

$\langle \Delta \vec{R}(\tau) \rangle$, mean drift

$\langle \Delta \vec{R}(\tau)^2 \rangle$, mean-squared displacement

τ , lag time

Definition: van Hove self correlation function

$$\mathbb{P}(\vec{r}, \tau) := \langle \delta(\vec{r} - \Delta \vec{R}(\tau)) \rangle \geq 0$$

probability of finding a displacement \vec{r}
after lag time τ

mean drift and mean-squared displacement
are moments of the van Hove function

$$\langle \Delta \vec{R}(\tau) \rangle = \int d\vec{r} \vec{r} \mathbb{P}(\vec{r}, \tau)$$

$$\langle \Delta \vec{R}(\tau)^2 \rangle = \int d\vec{r} r^2 \mathbb{P}(\vec{r}, \tau)$$

\mathbb{P} can be measured experimentally (particle

fracturing, microscopy...) and in
computer simulations

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Definition: **Incoherent (or Self-intermediate)
Scattering Function**

$$F^S(\vec{q}, \tau) = \langle \exp[-i\vec{q} \cdot \Delta \vec{R}(\tau)] \rangle \\ = \int d\vec{r} \hat{n}(\vec{r}, t) \exp[-i\vec{q} \cdot \vec{R}(\tau)]$$

$F^S(\vec{q}, \tau)$ is the characteristic function of
the random variable $\Delta \vec{R}(\tau)$

$F^S(\vec{q}, \tau)$ measured in incoherent scattering
experiment with photons (light, x-ray) or
particles (e.g. neutrons)

Definition: **Single-Particle Density**

$$g^S(\vec{r}, t) = \delta(\vec{r} - \vec{R}(t))$$

Fourier transform $g_q^S(t) = e^{i\vec{q} \cdot \vec{R}(t)}$

Incoherent scattering function is the
autocorrelation function of the single-
particle density

$$F^S(\vec{q}, \tau) = \langle g_q^S(t) \cdot g_q^S(t_0) \rangle, \tau = t - t_0$$

arrays $\langle g_q^S(t) \rangle = 0$ for $\vec{q} \neq 0$

$$F_{q=0}^S(t) = 1, |F_q^S(t)| \leq 1$$

Definition: Incoherent Dynamic Structure Factor

$$F_q^s(z) = i \int_0^\infty dt e^{izt} F^s(\vec{q}, t), \quad z \in \mathbb{C}_+$$

$\text{Im} [F_q^s(z = \omega + i0)]$ measured in scattering experiments

Examples

1) drift / directed motion $\vec{R}(t) = \vec{R}_0 + \vec{v}t$

\vec{R}_0 is averaged out, consider only

$$\Delta \vec{r} = \vec{R}(t) = \vec{v}t \rightarrow P(\vec{r}, \tau) = \delta(\vec{r} - \vec{v}\tau)$$

$$F^s(\vec{q}, \tau) = \exp(-i\vec{q} \cdot \vec{v}\tau) \quad \text{complex!}$$

does not decay, too!

↳ particle with fixed velocity is not in equilibrium

2) Brownian motion

$\Delta R_\alpha(\tau)$ are independent Gaussian variables, i.e., a random walk with non-trivial cumulants

$$\langle \Delta R_\alpha(\tau) \rangle = 0, \quad \langle \Delta R_\alpha(\tau) \Delta R_\beta(\tau) \rangle = 2D\tau \delta_{\alpha\beta}$$

$$\alpha, \beta = x, y, z$$

$$P(\vec{r}, \tau) = \frac{1}{\sqrt{(4\pi D\tau)^d}} e^{-\frac{r^2}{4D\tau}}$$

$$F^s(q, \tau) = e^{-Dq^2\tau}$$

3) Brownian motion with drift

$$\vec{R}(t) = \vec{v}t + \zeta(t)$$

drift
form

Brownian noise term

$$P(\vec{r}, \tau) = \frac{1}{\sqrt{(4\pi D\tau)^d}} \exp[-(\vec{r} - \vec{v}\tau)^2 / (4D\tau)]$$

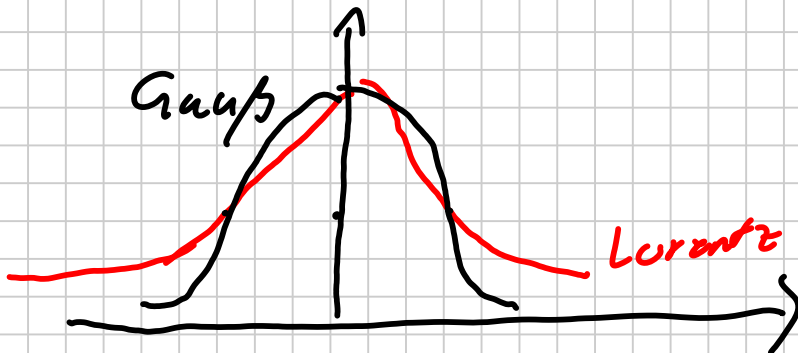
$$F^s(\vec{q}, \omega) = \exp[-i\vec{q} \cdot \vec{v}\tau - q^2 D\tau]$$

and
$$F_q^s(z) = \frac{-1}{z - \vec{q} \cdot \vec{v} + i D q^2}$$

$$\text{Im}[F_q^s(z)] = \frac{-\text{Im}[(z - \vec{q} \cdot \vec{v}) - i D q^2]}{(z - \vec{q} \cdot \vec{v})^2 + D^2 q^4} = \frac{D q^2}{(z - \vec{q} \cdot \vec{v})^2 + D^2 q^4}$$

Spectral line shape is a Lorentzian

with maximum at $\omega = \vec{v} \cdot \vec{q}$



4) general case

$$F^s(\vec{q}, \tau = t - t') = \langle \exp[-i\vec{q} \cdot (\vec{R}(t) - \vec{R}(t'))] \rangle$$

$$F^s(\vec{q}, \tau = 0) = 1$$

$$F^s(\vec{q}, \tau = 0) = -i\vec{q} \cdot \langle \dot{\vec{R}}(0) \rangle = i\vec{q} \cdot \vec{v}, \text{ drift}$$

usually not present

$$\begin{aligned} \ddot{F}^S(\vec{q}, \bar{v}=0) &= -\frac{d}{dt} \frac{d}{dt'} \bar{F}^S(\vec{q}, t-t') \Big|_{t=t'=0} \quad -157- \\ &= -\langle \vec{q} \cdot \ddot{\vec{R}}(t) q \cdot \ddot{\vec{R}}(t') \exp[-iq\{\vec{R}(t) - \vec{R}(t')\}] \rangle \end{aligned}$$

Definition: tagged current density

$$\vec{j}_q^S(t) = \ddot{\vec{R}}(t) \exp[iq \cdot \vec{R}(t)]$$

$$\begin{aligned} \ddot{F}^S(\vec{q}, t) &= -q_\alpha q_\beta \langle j_\alpha^S(\vec{q}, t) j_\beta^S(\vec{q}) \rangle = \\ &= -q^2 \langle j^{SL}(\vec{q}, t) j^{SL}(\vec{q}) \rangle \end{aligned}$$

current-current correlation function
for longitudinal self-carriers

Mori-Zwanzig projector $P = |S_q^S\rangle \langle S_q^S|$

particle conservation $\partial_t S_q^S(t) = i\mathcal{L} S_q^S(t) = iq j_q^{SL}(t)$

frequency matrix

$$\langle S_q^S | \mathcal{L} | S_q^S \rangle = \langle S_q^S | j_q^{SL} \rangle = \sigma, \text{ parity}$$

memory kernel

$$\begin{aligned} \langle \mathcal{A} \mathcal{L} S_q^S | \frac{1}{\mathcal{A} \mathcal{L} \mathcal{A} - z} | \mathcal{A} \mathcal{L} S_q^S \rangle = \\ = q^2 \langle j_q^{SL} | \frac{1}{\mathcal{A} \mathcal{L} \mathcal{A} - z} | j_q^{SL} \rangle \text{ as } \mathcal{A} | j_q^{SL} \rangle = | j_q^{SL} \rangle \end{aligned}$$

Definition *diffusion kernel*

$$\hat{D}(q, z) := \langle j_q^{SL} | \frac{1}{\mathcal{A} \mathcal{L} \mathcal{A} - z} | j_q^{SL} \rangle$$

→ general eqn

$$\hat{F}^S(q, z) = \frac{-1}{(z - \vec{q} \cdot \vec{v}) + q^2 \hat{D}(q, z)}$$

or in the time domain

$$\dot{F}^S(q, t) - i\vec{q} \cdot \vec{v} F^S(q, t) + \int_0^t dt' q^2 D(\vec{q}, t-t') F^S(q, t') = 0$$

Nb: approximately D is a safer alternative to approximately F^S directly, e.g.

$$\hat{D}(q, z) = (\vec{q} \cdot \vec{v})^2 \underset{\substack{| \\ \text{parallel to} \\ \text{drift}}}{D_{\parallel}} + [1 - (\vec{q} \cdot \vec{v})^2] \underset{\substack{| \\ \text{perpendicular} \\ \text{to drift}}}{D_{\perp}}$$

diffusion coefficient D from $\hat{D}(q, z)$:

velocity $v^S(t) = \lim_{q \rightarrow 0} j_q^S(t)$ in q direction

$$D = \lim_{z \rightarrow 0} \lim_{q \rightarrow 0} \frac{1}{z} \hat{D}(q, z) =$$

$$= \lim_{q \rightarrow 0} \int_0^{\infty} dt \langle v^S | e^{-iQzqt} | v^S \rangle$$

$$Q = 1 - P = 1 - |g_q^S\rangle \langle g_q^S|$$

$$\mathcal{L}Q = \mathcal{L} - \mathcal{L} |g_q^S\rangle \langle g_q^S| \rightarrow \mathcal{L} \text{ as } q \rightarrow 0$$

since g_q^S is conserved

$$\mathcal{L} \hookrightarrow \boxed{D = \int_0^{\infty} dt \langle v^S | e^{-i\mathcal{L}t} | v^S \rangle}$$

Green-Kubo relation